Empirical Estimation of Attenuation from Differential Propagation Phase Measurements at C Band

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(Manuscript received 24 November 2005, in final form 11 July 2006)

ABSTRACT

A polarimetric method is devised to correct for attenuation effects at C band on reflectivity $Z_H$ and differential reflectivity $Z_{DR}$ measurements. An operational cross-correlation analysis is used to derive advection vectors and to displace echoes over a 5-min time step. These advected echoes are then compared with observations valid at the same time. The method assumes that the mean change in the intrinsic $Z_H$ and $Z_{DR}$ over a 5-min period when considering 1–2 h of observations over the entire radar umbrella is approximately zero. Correction coefficients are retrieved through the minimization of a cost function that links observed decreases in $Z_H$ and $Z_{DR}$ due to attenuation effects with increases in differential phase shift ($\Phi_{DP}$). The retrieved coefficients are consistent with published values for the typical ranges of temperatures and drop sizes encountered at midlatitudes, even when Mie scattering effects are present. Measurements of $Z_H$ and $Z_{DR}$ corrected using retrieved coefficients are compared with raw measurements and to measurements adjusted by mean coefficients found in the literature. The empirical retrieval method shows improvement over using mean correction coefficients based on comparisons of $Z_H$ from neighboring, unattenuated radars, disdrometer measurements, and analysis of $Z_H$ and $Z_{DR}$ as a function of $\Phi_{DP}$.

1. Introduction

Improved sensitivity and relatively lower costs of radars operating at C-band frequencies may be partially offset by their susceptibility to attenuation that can significantly reduce values of reflectivity $Z_H$ and differential reflectivity $Z_{DR}$, thus affecting the quality of rainfall measurements. Attenuation correction methods that rely on $Z_H$ measurements alone have proven to be unreliable. Hitschfeld and Bordan (1954) proposed a scheme based on an empirical relationship between the specific horizontal attenuation $A_H$ and $Z_H$. Although Aydin et al. (1989) and Gorgucci et al. (1995, 1998) suggested accuracies of gate-to-gate corrections up to 10%, these profiling schemes are notoriously unstable because any initial error in $Z_H$ due to a slight miscalibration grows very rapidly.

Bringi and Chandrasekar (2001) provide an excellent summary of methods to correct for attenuation of $Z_H$ and $Z_{DR}$ at attenuating frequencies using the phase difference $\Phi_{DP}$ between horizontally ($H$) and vertically ($V$) polarized returns. These measurements have the advantage of being immune to radar miscalibration (Zrnić and Ryzhkov 1996). As raindrops grow they become more oblate, leading to positive values of $Z_{DR}$, and the slower velocity of the $H$ wave relative to $V$ leads to a value of $\Phi_{DP}$ that increases with range in a rainfall medium. Scattering simulations assuming gamma functions for raindrop size distributions suggested that the attenuation $A_H$ and differential attenuation $A_{HV}$ are nearly linearly related to the gradient of $\Phi_{DP}$ with range in degrees per kilometer or specific differential phase $K_{DP}$ at X, C, and S bands by coefficients $a$ and $b$, respectively (Bringi et al. 1990). These coefficients can vary because of changes in drop tem-

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DOI: 10.1175/JAM2464.1

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Ryzhkov and Zrnić (1995) produced scatterplots of $\Phi_{DP}$ versus $Z_H$ and $\Phi_{DP}$ versus $Z_{DR}$ using a large dataset collected at S band. The slope of the fitted lines yielded values of $a$ and $b$ that were larger than those of Bringi et al. (1990), which they attributed to additional attenuation due to Mie scattering. Carey et al. (2000) adapted the technique of Ryzhkov and Zrnić (1995) for C-band radar by identifying rays with Mie scattering effects (equivolumetric median diameter drops $> 2.5$ mm or $Z_{DR} > 2.5$–3 dB at C band) where larger coefficients (2 times as large for $a$ and 4 times as large for $b$!) were applied to these rays if they had large backscatter differential phase $\delta$, dips in the copolar cross-correlation coefficient at zero time lag $[\rho_{HH}(0)]$, and high values of $K_{DP}$. Ranges of $K_{DP}$ are also used to limit the data sample so that the intrinsic scatter in $Z_{DR}$ and $Z_H$ at a given value of $\Phi_{DP}$ is minimized.

Smyth and Illingworth (1998) proposed a method at C band based on constraining values of $Z_{DR}$ behind intense convective cells where the intrinsic $Z_{DR}$ was believed to be 0 dB, representing spherical droplets. Observations of negative $Z_{DR}$ were used to correct for horizontal and differential attenuation effects; the total differential attenuation was then redistributed to gates along the radial with $K_{DP} > 1$° km$^{-1}$, yielding an estimate of $A_{HV}$. A linear relationship was then assumed between $A_{HV}$ and $A_k$ based on scattering simulations, thus providing for a correction to values of $Z_H$ and $Z_{DR}$. This method offers the advantage of being applied on a ray-by-ray basis. However, the constraint of an intrinsic $Z_{DR}$ of 0 dB in drizzle regions, for which the method is based, is difficult to identify automatically using other polarimetric parameters. In addition, the approach is susceptible to nonuniform beam-filling effects on $Z_{DR}$ measurements near the edges of convection.

The “ZPHI” algorithm (Testud et al. 2000) is similar in design to the original technique of Hitschfeld and Bordan (1954), with the additional constraint of estimating the reference attenuation at C band. The $\Phi_{DP}$ constraint leads to numerical stability, which was the main failing with the original technique. The ZPHI algorithm may yield inaccurate results because of deviations from the assumed raindrop size–shape model and in situations where the observed $\Phi_{DP}$ is small relative to system noise levels. Also, $Z_H$ values become biased along with the estimated $A_H$ if measurements from hail or mixed-phase hydrometeors are included in the calculations. This method has been modified by adding a combined $\Phi_{DP}$-$Z_{DR}$ constraint as described in Bringi and Chandrasekar (2001, their section 7.4) and has been adapted for use at X band (Iwanami et al. 2003; Anagnostou et al. 2004; Park et al. 2005a).

A new approach is proposed herein that seeks a solution to the coefficients $a$ and $b$ linking attenuation of $Z_H$ and $Z_{DR}$ to observations of $\Phi_{DP}$. The technique is empirical and is based on the assumption that the mean, intrinsic difference between cross-correlated $Z_H$ ($Z_{DR}$) data over a 5-min period is approximately zero when a sufficiently large number of observations are considered. Reductions in $Z_H$ ($Z_{DR}$) are shown to be linearly related to increases in $\Phi_{DP}$, which are used to solve for the $a$ and $b$ coefficients. Coefficients are retrieved for seven cases of intense convection occurring near Météo-France’s C-band polarimetric radar. Evaluations are made possible by comparing raw measurements with measurements adjusted using mean coefficients found in the literature, and to corrected $Z_H$ and $Z_{DR}$ values using empirically retrieved coefficients.

The conceptual basis and mathematical background of the technique are described in section 2. Section 3 presents results of the retrieved curves. The variability of these curves from case to case is also examined in this section. Section 4 evaluates the accuracy of the method by examining the behavior of $Z_H$ and $Z_{DR}$ as a function of $\Phi_{DP}$ with uncorrected data, data corrected using mean coefficients, and data corrected using retrieved coefficients. The methods are also objectively evaluated using comparisons with unattenuated data from neighboring radars and disdrometer measurements. A summary of results and conclusions are presented in section 5. The final section also discusses operational aspects of the algorithm, including its limitations and potential improvements.

2. Description of method to retrieve attenuation and differential attenuation coefficients to correct radar signals at C band

a. Conceptual overview

The retrieval technique was envisioned based on observations of $Z_{DR}$, and to a lesser extent $Z_H$, as intense convective cells moved near the Trappes, France, radar. Strong rain cells at close range were noted to have the effect of producing wedge-shaped artifacts of anomalously low $Z_{DR}$ values at ranges beyond the initial, attenuating cell. These wedges would often propagate in the azimuthal direction with respect to the radar, creating the appearance of a moving searchlight. Images of $Z_H$ also revealed these artifacts, yet they were harder to
discern 1) because of the higher variability of $Z_H$ in attenuated regions as compared with $Z_{DR}$ and 2) because attenuation reduces $Z_{DR}$ values more substantially per unit $Z_{DR}$ as compared with $Z_H$. Because these artifacts were moving quickly from scan to scan, it occurred that it would be possible to advect $Z_H$ ($Z_{DR}$) data from a previous scan so that the data were matched in space in time. Mean differences are assumed to be due to attenuation and are thus linked to differences in $\Phi_{DP}$ measurements between advected and observed images.

Figure 1 illustrates the empirical retrieval concept developed herein using two scans of $Z_{DR}$ and $\Phi_{DP}$ at an elevation angle of 1.5°, where the data from the first scan have been advected using an operational cross-correlation analysis so they match the data in the second scan. A wedge of negative $Z_{DR}$ values is evident northwest of the radar (Fig. 1a) that was the result of an
intense convective cell very near the radar. After 5 min the cell moved north of the radar so that the wedge of attenuated $Z_{DR}$ data moved clockwise with respect to the radar, resembling a moving searchlight (Fig. 1b). Figure 1c shows that this wedge of attenuated $Z_{DR}$ was coincident with large values of $\Phi_{DP}$. Large $\Phi_{DP}$ values also rotated in a searchlight pattern to the north of the radar (Fig. 1d). If we assume for simplicity that the $Z_{DR}$ data to the northwest of the radar were from spherically shaped drops that changed from being attenuated (−5 dB) to unattenuated (0 dB), then these increases can be linked to reductions in $\Phi_{DP}$ from 100° to 0°, for example. In this case, we can say that a gain of 5 dB in $Z_{DR}$ was caused by a loss of 100° in $\Phi_{DP}$, so that the $b$ coefficient is 0.02. However, such a direct solution at a particular grid point is subject to errors caused by intrinsic, microphysical changes within a 5-min period.

Thousands of comparisons are available using each $Z_{DR}$, $Z_{DR}$ pair from a series of radar scans. A data pair refers to the advected and observed values of $Z_{DR}$ and $Z_{DR}$ at a given grid point. Mean differences in the intrinsic $Z_{DR}$ and $Z_{DR}$ between the advected and observed images are assumed to be zero when considering thousands of data pairs under the radar umbrella over a 1–2-h duration. Changes in a 5-min period do occur at a given pixel because of precipitation growth, decay, or microphysical processes. However, the net change over the entire spatial/temporal domain is assumed to be zero. In other words, it is equally probable that a given precipitation feature will change intrinsically because of growth or decay, and the net change when considering thousands of pairs is thus zero.

Raw measurements of $Z_{DR}$, $Z_{DR}$, copolar cross-correlation coefficient at zero time lag $\rho_{HV}(0)$, and $\Phi_{DP}$ from the Trappes radar are subject to errors resulting from mismatched data, near-radome interference effects, and noise. For details regarding an examination of each of these effects, proposed correction techniques, and a summary of the data quality of the Trappes polarimetric radar, see Gourley et al. (2006). In summary, structures in the near field of the antenna were found to affect $Z_{DR}$ measurements by as much as 0.4 dB. An empirical mask was developed and subsequently applied to all $Z_{DR}$ data used in this study. Biases in $Z_{DR}$ and $\rho_{HV}(0)$ measurements that occur at signal-to-noise ratios of less than 10 dB have been mitigated. Measurements of $Z_{DR}$ have been calibrated using observations at vertical incidence where the intrinsic $Z_{DR}$ is known to be 0 dB. The behavior of initial differential phase measurements with the Trappes radar has been examined, and it was found that raw measurements are aliased, vary with azimuth, and have a mean offset of 6°. An empirical equation was developed to correct $\Phi_{DP}$ measurements for the initial system offset and has subsequently been applied to all $\Phi_{DP}$ measurements used in this study.

Thresholds are applied to the data pairs so that errors from brightband contamination, ground clutter, hail contamination, and noisy $\Phi_{DP}$ measurements in light rain are mitigated. Values of $\rho_{HV}(0)$ were observed to decrease below 0.97 within the melting layer. Thus, data pairs are discarded from the retrieval method if they have values of $\rho_{HV}(0)$ less than 0.97. This threshold also eliminates many $Z_{DR}$ and $Z_{DR}$ measurements affected by ground clutter. Data must be collected within 5–150 km from radar to reduce ground clutter and brightband contamination further. Measurements in hail are avoided by thresholding maximum $Z_{DR}$ values at 40 dBZ. Noisy $\Phi_{DP}$ measurements are first mitigated by smoothing $\Phi_{DP}$ profiles along a 25-gate window corresponding to 6 km. Second, the difference in $\Phi_{DP}$ between advected and observed data pairs must be more than 10° so that the comparisons are made well above the measurement noise level of 1.8° found with the Trappes radar (Gourley et al. 2006).

The developed method is dependent on motion vectors derived from a cross-correlation analysis that is used to displace $Z_{DR}$, $Z_{DR}$, and $\Phi_{DP}$ data forward in time by 5 min. The cross-correlation technique is very similar to the Tracking Radar (or Reflectivity) Echoes by Correlation (TREC) developed by Rinehart (1979) and later modified by Tuttle and Foote (1990). An array of reflectivity data within a 60 × 60 km² area is correlated with a second array of reflectivity data separated by 5 min. The displacement between the initial array location and that of the array having the largest correlation coefficient determines the motion vector. These derived motion vectors are then used to advect the echoes forward in time by 5 min. Errors in the motion vectors will cause the data initially separated by 5 min to be mismatched. In this case, differences in $Z_{DR}$ ($Z_{DR}$) for a given difference in $\Phi_{DP}$ will not be due to attenuation but rather to displacement. The effect on the retrieved coefficients depends on the structure of the precipitation relative to the biases in the motion vectors, but in general the retrieved coefficients are expected to be normally distributed around a mean value of zero.

It should be noted that the retrieval technique does not require that one of the measurements remain completely free from attenuation effects. It is the mean difference in $Z_{DR}$ ($Z_{DR}$) between measurements that have been advected over a 5-min period and observed data that is of interest to the algorithm. A relationship is assumed between domainwide $Z_{DR}$ ($Z_{DR}$) differences per unit change in $\Phi_{DP}$, thus providing for the computation of the $a$ ($b$) coefficient.
b. Mathematical background

The goal of this analysis is to derive empirically the coefficients \( a \) and \( b \) that link the observed \( \Phi_{Dp} \) to the attenuation of \( Z_H \) and \( Z_{DR} \). For simplicity, the development presented here is for \( Z_H \) and is completely analogous to the treatment of \( Z_{DR} \). It is first assumed that the measured reflectivity \( Z_H \) is related to the intrinsic reflectivity as follows:

\[
Z_H = Z_H^{\text{int}} - a \Phi_{Dp}. \tag{1}
\]

where \( a \) is the coefficient for horizontal attenuation [dB (°)^{-1}]. Because (1) relies on an absolute measure of \( \Phi_{Dp} \), measurements have been dealiased and corrected for the initial system offset as discussed in section 2a.

Next, we assume the relationship in (1) holds for

\[
Z_{H}(i) - \hat{Z}_{H}(i) = (-1, 0, 0, 0, 0, 0, 0, 1, \ldots) \cdot \left( \begin{array}{c}
A(0) \\
A(1) \\
\vdots \\
A(360)
\end{array} \right), \tag{4}
\]

where, in the example shown in (4), \( \Phi_{Dp} \) values occupy the 1st and 10th position element, meaning the observed and advected \( \Phi_{Dp} \) values were 0° and 10°, respectively. The 360 elements correspond to all possible \( \Phi_{Dp} \) values ranging from 0° to 360°. This enables the total attenuation \( A \) to be computed as a function of \( \Phi_{Dp} \). This calculation is repeated over all \( n \) measurements corresponding to all ranges and azimuths for a 1–2-h duration. Equation (4) is more easily represented in the following matrix system:

\[
Z - \hat{Z} = \Phi_{Dp} \cdot (-A), \tag{5}
\]

where observed and advected reflectivity vectors \( Z \) contain \( n \) elements, \( \Phi_{Dp} \) is an \( n \times 360 \) matrix, and \( A \) is a 360-element vector corresponding to the coefficients. Because the system is overdetermined (\( n \gg 360 \)), the following cost function \( J \) is introduced:

\[
J = [Z - \hat{Z} - \Phi_{Dp} \cdot (-A)]^T \cdot [Z - \hat{Z} - \Phi_{Dp} \cdot (-A)]
+ \lambda \left[ U^T \cdot (-A) \right] \cdot [U^T \cdot (-A)], \tag{6}
\]

where the superscript \( T \) refers to the transpose of the matrix. To make the retrieval well posed, a term involving \( U \) is added to the cost function imposing the first element of \( A \) to be 0 (i.e., there is no attenuation if \( \Phi_{Dp} \) is equal to zero). The vector \( U \) is composed of 360 elements, where the first is equal to 1 and the remaining ones are set to 0. The weight of this constraint is controlled by setting a very large value for \( \lambda \). The cost function \( J \) is minimized with respect to \( A \) in a least squares sense as follows:

\[
A = -[\Phi_{Dp}^T \cdot \Phi_{Dp} - \lambda U^T \cdot U]^{-1} \cdot \Phi_{Dp}^T \cdot (Z - \hat{Z}). \tag{7}
\]

This method provides for the retrieval of the total attenuation (total differential attenuation) in decibels at all 360 values of \( \Phi_{Dp} \). If linearity is assumed, then slopes of lines fit to the retrieved path-integrated attenuation (differential attenuation) as a function of \( \Phi_{Dp} \) correspond to the coefficients \( a \) (b) used to correct observations of \( Z_H (Z_{DR}) \). The following section examines this linearity assumption and the variability of retrieved coefficients with \( \Phi_{Dp} \).

3. Results from attenuation and differential attenuation retrievals

a. Case studies

The technique of retrieving coefficients for attenuation using advected fields of \( Z_H \) and \( Z_{DR} \) is tested here for seven cases of intense convection. As seen in Table 1, a very large sample size is produced using as few as 12 scans (one hour’s worth) at an elevation angle of 1.5°. Figures 2a,b show the path-integrated attenuation and differential attenuation plotted as a function of
Linear least squares regression is then used to fit a line to each curve and obtain the slope; these slopes are the retrieved coefficients and are summarized in Table 1. The thick gray lines in Figs. 2a,b correspond to the expected minimum and maximum correction coefficients at C band as reported in Carey et al. (2000). The minimum values come from scattering simulations, and the maxima are the empirical coefficients found by Carey et al. (2000) that account for increasing slopes resulting from the presence of large-diameter drops in convective cores.

The curves that were retrieved for attenuation of \( Z_H \) (Fig. 2a) nearly fall within the expected bounds. The 4 July 2005 case suggests there is the least amount of attenuation, and the 3 and 23 June 2005 cases indicate the most (refer to Table 1). The variability of these curves from case to case is examined in the next section. All curves describing the path-integrated differential attenuation (Fig. 2b) fall within expected bounds. These curves are better approximated as being linear and are smoother than those shown in Fig. 2a. The following section examines whether the different \( a \) and \( b \) coefficients that were retrieved are related to the presence of large drops or different temperatures in the attenuating cells.

### Variability of retrieved correction coefficients

Attenuation effects on \( Z_H \) and \( Z_{DR} \) enhance as equivolume median diameters \( D_o \) enter the Mie scattering regime at approximately 2.5 mm at C band (Carey et al. 2000). Simulations have also shown that the temperature of the raindrops, DSD variability, and raindrop-axis-ratio relation affect the relationship between attenuation and \( \Phi_{DP} \) (Jameson 1992; Matrosov et al. 2002, 2005). The Carey et al. method improves over the original technique of Ryzhkov and Zrnić (1995) by accounting for the variability of \( a \) and \( b \) due to large drops resulting in Mie scattering. Although

**Table 1. Summary of event dates and times, sample sizes, and retrieved attenuation correction coefficients [dB (°)\(^{-1}\)]**

<table>
<thead>
<tr>
<th>Date</th>
<th>Event start-end times (UTC)</th>
<th>Sample size</th>
<th>( a )</th>
<th>( b )</th>
</tr>
</thead>
<tbody>
<tr>
<td>24 Mar 2005</td>
<td>1600–1800</td>
<td>6.95 \times 10^6</td>
<td>0.0805</td>
<td>0.0304</td>
</tr>
<tr>
<td>03 Jun 2005</td>
<td>1400–1600</td>
<td>7.97 \times 10^4</td>
<td>0.1111</td>
<td>0.0311</td>
</tr>
<tr>
<td>23 Jun 2005</td>
<td>1500–1700</td>
<td>1.62 \times 10^6</td>
<td>0.1044</td>
<td>0.0496</td>
</tr>
<tr>
<td>26 Jun 2005</td>
<td>1000–1100</td>
<td>1.01 \times 10^6</td>
<td>0.0817</td>
<td>0.0268</td>
</tr>
<tr>
<td>28 Jun 2005</td>
<td>2000–2100</td>
<td>1.04 \times 10^6</td>
<td>0.0719</td>
<td>0.0294</td>
</tr>
<tr>
<td>30 Jun 2005</td>
<td>1500–1700</td>
<td>1.61 \times 10^6</td>
<td>0.0737</td>
<td>0.0336</td>
</tr>
<tr>
<td>04 Jul 2005</td>
<td>0400–0500</td>
<td>1.76 \times 10^6</td>
<td>0.0300</td>
<td>0.0166</td>
</tr>
</tbody>
</table>

**Fig. 2.** Retrieved curves (in black) of \( \Phi_{DP} \) vs (a) path-integrated attenuation and (b) path-integrated differential attenuation for seven cases of intense convection listed in Table 1. The slopes of the gray lines correspond to the minimum and maximum correction coefficients published in simulation-based and empirical studies.
simulations indicate that $a$ and $b$ are sensitive to the assumed drop shape model, Matrosov et al. (2005) note that the slope of the drop-axis-ratio relation is fairly constant except for drops with diameters larger than 7 mm. The retrieval method developed herein is empirical; therefore, separating these effects is not unambiguous.

The $Z_{\text{DR}}$ polarimetric variable responds well to hydrometeor shapes, sizes, and habits, and therefore it can be used to detect the presence of large drops ($3 < Z_{\text{DR}} < 5$ dB at C band) provided that the $Z_{\text{DR}}$ data are not attenuated. To avoid attenuated $Z_{\text{DR}}$ values, a threshold of $\Phi_{\text{DP}} < 10^\circ$ is implemented so that the analysis is effectively restricted to portions of cells that are either close to the radar or have very little precipitation between them and the radar. The percent of unattenuated data having $Z_{\text{DR}}$ values within the Mie scattering regime is reported. The temperatures of these cells are also determined by utilizing radiosonde observations launched at the Trappes radar site.

Figure 3a shows the retrieved $a$ and $b$ coefficients plotted against the percent of unattenuated cells containing big drops for all cases listed in Table 1. Although the data sample is small, both coefficients exhibit a nonlinear dependence on the presence of big drops within these cells. Increases in $a$ and $b$ are most significant up to percentages of 10% and level off thereafter. The drop size dependence of relationships between attenuation and $\Phi_{\text{DP}}$ cannot be neglected. The processes resulting in high sensitivity of $a$ and $b$ to drop sizes are either Mie scattering effects or changes in the drop-axis-ratio relation. Regardless of the mechanism, the empirical technique developed herein is capable of retrieving coefficients to account for the presence of big drops. If attenuating cells with big drops are present, then larger differences in $Z_H$ ($Z_{\text{DR}}$) over a 5-min period at farther regions will be observed for a given value of $\Phi_{\text{DP}}$. Larger $a$ ($b$) coefficients will be derived and then may be used subsequently to correct values of $Z_H$ ($Z_{\text{DR}}$). It is noted, however, that the retrieved coefficients in these regions may have the effect of overcorrecting $Z_H$ and $Z_{\text{DR}}$ data that do not contain big drops. Last, there is a linear correlation between the $a$ and $b$ coefficients from case to case of 0.75. These are derived independently, and empirical evidence presented here confirms prior correlations discovered through simulations.

The dependence of retrieved coefficients on temperature is shown in Fig. 3b. Coefficients would adapt accordingly in the developed scheme if there were indeed a strong sensitivity to temperature. Similar to the scenario presented above with big drops, larger changes in $Z_H$ ($Z_{\text{DR}}$) for a given value of $\Phi_{\text{DP}}$ would be observed over a 5-min period behind regions of anomalously cold raindrops. The retrieved coefficients would therefore be greater at colder temperatures. However, this effect is not evident in Fig. 3b. It is possible that the temperature dependence of coefficients has been overshadowed by the strong dependence on drop sizes.

4. Validation of retrieved attenuation and differential attenuation coefficients at C band

a. Variables evaluated as a function of $\Phi_{\text{DP}}$

If we are to assume that the distribution of $Z_H$ and $Z_{\text{DR}}$ is independent of range from radar location and
there is no attenuation, then both variables should also be independent of $\Phi_{DP}$ measurements. The validity of this assumption increases with stratiform precipitation, which is more spatially homogeneous than in convection, and for a long duration of analysis. Analyzing corrected and uncorrected $Z_{DR}$ and $Z_H$ as a function of $\Phi_{DP}$ serves as a tool for evaluating the accuracy of the retrieved coefficients, whereas it is the basis of deriving $a$ and $b$ coefficients in the Carey et al. (2000) method. Matrosov et al. (2005) adopt a similar strategy but fix their analysis to a ground-clutter pixel where the reflectivity is assumed to be relatively constant.

Figure 4 shows observations of $Z_H$ and $Z_{DR}$ with no correction for attenuation, correction based on a fixed set of coefficients, and correction based on empirically retrieved coefficients plotted as a function of $\Phi_{DP}$ for the 23 June 2005 case. The fixed parameters $a = 0.0688$ dB $(^\circ)^{-1}$ and $b = 0.01785$ dB $(^\circ)^{-1}$ come from mean values at C band based on a literature survey reported in Carey et al. (2000). The 23 June 2005 case is of particular interest because the retrieved $a$ ($b$) coefficient is 52% (178%) larger than the literature-mean value.

Figure 4a shows the uncorrected $Z_H$ data have a clear decreasing trend with increasing $\Phi_{DP}$ that is due to attenuation effects. Reflectivity data corrected using a fixed $a$ coefficient have no evident bias as a function of $\Phi_{DP}$. The retrieved $a$ coefficient apparently overcorrects $Z_H$ and results in errors as large as 5 dB. It is possible that the intrinsic $Z_H$ increases with $\Phi_{DP}$ because the analysis covers only a 2-h duration. However, it is more likely that errors in the developed retrieval method result from a single coefficient being derived and then applied to gates at several ranges and azimuths. The large coefficients retrieved from the method may only be applicable to the regions behind the attenuating cells containing large drops. The application of a mean coefficient to the entire field of observations is a limitation of the retrieval method in its current configuration.

Observations of uncorrected $Z_{DR}$ data for the same event show a similar decreasing trend with increasing values of $\Phi_{DP}$ (Fig. 4b). Mean values of $Z_{DR}$ approach $-7$ dB at $\Phi_{DP}$ measurements of 160°. In this case, corrections using the literature-mean $b$ coefficient result in $Z_{DR}$ values that are biased too low by at least 5 dB for $\Phi_{DP}$ measurements of 160°. The fact that the literature-mean $a$ coefficient provides an accurate correction but the $b$ coefficient does not suggests that the literature-mean coefficients are not consistent to each other. After $Z_{DR}$ observations have been adjusted using the empirically retrieved $b$ coefficient, the data show much less dependence on $\Phi_{DP}$. There is a slight reduction with increasing $\Phi_{DP}$, but it is evident the retrieved $b$ offers a significant improvement over the literature-mean coefficient.

b. Radar–disdrometer comparisons

Similar to the analysis employed in Park et al. (2005b), a disdrometer is used to evaluate the accuracy of the attenuation and differential attenuation correction schemes. A laser-optical disdrometer was placed 24 km from the Trappes radar along the 77° azimuth in Parc Montsouris, Paris. The laser precipitation monitor produces a 228-mm-long parallel infrared beam that is received by a photodiode with a lens. The amplitude of the reduction of the received signal is used to calculate the diameter of the particles passing through the beam. Drop spectra are measured every minute by the disdrometer for drops ranging from 0.1875 to 7 mm in diameter. Reflectivity at horizontal and vertical polarization is computed using scattering simulations at C
band. Drop temperatures are assumed to be 15°C and the drop shape model from Goddard et al. (1995) is used. Polarimetric variables computed from disdrometer measurements are subject to errors from the assumed temperature and drop shape model. Another significant source of uncertainty in this analysis comes from the vastly different scales at which disdrometers and radars measure polarimetric quantities. Regardless, simulated values from disdrometer measurements offer an independent measure of $Z_H$ and $Z_{DR}$.

Figure 5 shows a 2-h time series of radar-measured $Z_H$ and $Z_{DR}$ over the disdrometer that are uncorrected, corrected using fixed parameters, and corrected using retrieved parameters. At 1513 UTC, $\Phi_{DP}$ reached 177° over the disdrometer location and subsequently resulted in an attenuated $Z_H$ that was at least 10 dB below simulated values from the disdrometer (Fig. 5a). When the $Z_H$ data have been corrected using a fixed coefficient, measurements agree with simulated values from the disdrometer more closely. However, the best agreement comes from $Z_H$ data that have been corrected with the larger value retrieved using the empirical technique. Figure 4a suggests the retrieved $a$ coefficient for the 23 June 2005 case did not improve $Z_H$ measurements over using a fixed coefficient when we consider all data under the radar umbrella. Figure 5a, on the other hand, shows that the retrieved coefficient is more accurate than the fixed value when the analysis is limited to a particular bin that undergoes severe attenuation.

The uncorrected radar measurements of $Z_{DR}$ are attenuated by approximately 9 dB near 1513 UTC (Fig. 5b). Values of $Z_{DR}$ corrected using a fixed $b$ value underestimate disdrometer-based measurements by 6 dB during this period of severe differential attenuation. There is almost no bias in comparing disdrometer-based $Z_{DR}$ values with those corrected using the empirically retrieved $b$ coefficient. Independent measurements of polarimetric variables afforded by a disdrometer indicate the retrieved coefficients are more accurate than mean coefficients found in the literature for a case of severe attenuation.

c. Radar–radar reflectivity comparisons

There are three, nonpolarimetric C-band radars in the vicinity (Abbeville, Bourges, and Falaise) that operate continuously as part of the French radar network. Measurements from these neighboring radars are used to evaluate objectively the accuracy of retrieved and literature-mean coefficients used to correct $Z_H$. A similar approach could be implemented for evaluating retrieved $b$ coefficients provided the neighboring radars are polarimetric. Reflectivity values from adjacent radars are compared if $\rho_{HV}(0)$ is greater than 0.97 from the Trappes radar, difference in height of measurement is less than 500 m, range to echoes is less than 150 km from both radars, and attenuation from the nonpolarimetric radar is less than 1 dB. The latter criterion is based on a path-integrated $Z_H$ estimation from Doviak and Zrnić (1993).

Reflectivity comparisons are made every 5 min from 0000 to 2345 UTC 24 March 2005, resulting in $1.8 \times 10^5$, $5.1 \times 10^4$, and $2.6 \times 10^5$ data pairs for Trappes and Abbeville, Bourges, and Falaise, respectively. At least 50% of these data pairs were collected at values of $\rho_{HV}$ in all cases. Nonzero $Z_H$ differences were noted even at low $\rho_{HV}$ values ($<10^5$) for which attenuation effects can be neglected. Gourley et al. (2003) and Tabary (2003) showed that the temporal variability of a radar’s calibration can be monitored by comparing $Z_H$ from two different radars. Differences that are not due to attenuation were found by computing the average $Z_H$ difference from Trappes and each neighboring nonpolarimetric C-band radar for $0^6 < \Phi_{DP} < 10^6$ over the duration of the event. Reflectivity differences
caused by the combined effects of radar calibration differences and variable beam propagation paths were found to be 1.9, 2.8, and −0.2 dB for Trappes minus Abbeville, Bourges, and Falaise, respectively, and were subsequently subtracted out from all reflectivity differences.

Figure 6 shows uncorrected $Z_H$ values as measured by Trappes minus $Z_H$ from neighboring radars are reduced as $\Phi_{DP}$ (a proxy for attenuation) increases. A literature mean coefficient of 0.0688 dB (°)$^{-1}$ and an empirically retrieved coefficient of 0.0805 dB (°)$^{-1}$ were then applied to each $Z_H$ measurement from the Trappes radar and subsequently were compared to all three neighboring radars. In all cases, the mean reflectivity difference is closest to 0 dB when the empirically retrieved coefficient is used to correct for attenuation. It is noted that the estimation of the mean becomes more uncertain as the sample size is reduced. This is especially noticeable at the largest values of $\Phi_{DP}$ where the curves exhibit sudden changes.

5. Summary and conclusions

An empirical method has been devised to retrieve coefficients to correct for power losses at C band due to attenuation and differential attenuation in rain. The method relies on the assumption that mean intrinsic values of $Z_H$ and $Z_{DR}$ for each cross-correlated precipitation feature under the radar umbrella do not change when at least $10^4$ samples over a 1–2-h duration are considered. Differences between advected and observed values of $Z_H$ and $Z_{DR}$ are linked to changes in $\Phi_{DP}$ and enable the empirical retrieval of $a$ and $b$ attenuation correction coefficients. The retrieved coefficients fall within expectations found through simulations and in other experimental results. Also, the linearity assumption between attenuation and $\Phi_{DP}$ was not initially assumed in this method but was found to be valid.

Relationships between retrieved coefficients and characteristics of cells that were responsible for the attenuation were analyzed. It was discovered that larger...
coefficients for both $Z_H$ and $Z_{DR}$ were empirically retrieved when a relatively large percentage of big drops ($Z_{DR} > 3$ dB at C band) were present. Moreover, the $a$ and $b$ coefficients were found to be linearly related with a correlation of 0.75. Temperature effects on attenuation, however, were not noticeable but were perhaps masked by the large dependence on the presence of large drops. Changes in raindrop shape or Mie scattering effects are likely culprits for the increased attenuation observed with large-diameter drops.

The accuracy of the retrieved coefficients was evaluated by plotting $Z_H$ and $Z_{DR}$ before and after correction as a function of $\Phi_{DP}$, comparison with disdrometer-based measurements, and comparison with unattenuated data from neighboring radars. Because improvements over uncorrected values were inevitable, a benchmark was established by correcting $Z_H$ and $Z_{DR}$ data using fixed $a$ and $b$ coefficients from a literature survey at C band. All analyses indicate the empirically retrieved $b$ coefficients result in more accurate $Z_{DR}$ values as compared with uncorrected values and those corrected using a fixed coefficient. There is a suggestion that the retrieved $a$ coefficient overcorrected $Z_H$ data when considering all observations around the radar. However, when a single pixel over the disdrometer site was examined, better agreement with simulated $Z_H$ from the disdrometer observations was accomplished using the empirically retrieved coefficient. This pixel had an associated $\Phi_{DP}$ value of $177^\circ$ and therefore was severely attenuated. It is plausible that the retrieval method responds to severe attenuation caused by Mie scattering and thus may overcorrect $Z_H$ observations that did not experience these effects.

Several considerations must be made regarding the operational application of this empirical technique. As with other attenuation correction methods, there is no consideration for inflated $Z_H$ values resulting from hail or brightband contamination. These errors will likely grow after attenuation coefficients have been applied. Also, it is noted that the method in its current form does not provide the capability of varying the coefficients in space. Instead, coefficients are applied to entire fields of $Z_H$ and $Z_{DR}$ and thus may overcorrect data that do not contain large drops, for example. Future methods should consider subdividing radar observations that have similar characteristics and then retrieving multiple coefficients. Approximately 1–2 h of data were needed to retrieve the coefficients successfully. These hours were typically characteristic of large wedges of $\Phi_{DP}$ that rotated rapidly in the azimuthal direction, resembling a searchlight. Data pairs are considered in the retrieval method if their values of $\Phi_{DP}$ differ by at least $10^\circ$. It will be more difficult to meet this criterion in stratiform rain that has a more homogeneous spatial structure.

Acknowledgments. This work was done in the frame of the Projet Application Radar à la Météorologie Infra-Synoptique (ARAMIS) Nouvelles Technologies en Hydrométéorologie Extension et Renouvellement (PANTHERE) supported by Météo-France, the French “Ministèr e de L’Écologie et du Développement Durable,” the “European Regional Development Fund (ERDF)” of the European Union, and CEMAGREF. The authors would like to acknowledge the advice provided by the PANTHERE scientific review committee. Specifically, Anthony Illingworth provided us with many useful insights regarding attenuation correction methods, and he helped with the review of this paper. Kim Dokhae was helpful in processing the disdrometer data. The comments from three anonymous reviewers greatly improved the quality of this manuscript. Their assistance is greatly appreciated.

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